

# Random spheres as a 1+1 dimensional field theory.

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**Abstract:** We construct non-gaussian Brownian spheres on  $R^d$  as a 1+1 dimensional field theory.

**Keywords:** Brownian bridge; infinite dimensional Brownian motion.

**MSC 2000 codes:** 60G60. 81T40. 82B31.

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# 1 Introduction

In string theory or in conformal field theory (See survey of Witten [51] or Gawedzki [18], [19],[20]), people look at a Riemannian surface  $\Sigma$  (the world-sheet) and at a Riemannian manifold  $M$  (the target space) and consider the set of maps  $x(\cdot)$  from  $\Sigma$  into  $M$  endowed with the formal probability measure:

$$(1) \quad d\mu(x(\cdot)) = Z^{-1} \exp[-I(x(\cdot))] dD(x(\cdot))$$

where  $I(x(\cdot))$  is the energy of the field  $x(\cdot)$ ,  $dD(x(\cdot))$  the formal Lebesgue measure on the set of fields and  $Z$  the partition function destined to get a probability measure.

If we consider the case where  $M = R$  with the flat Riemannian structure, we get a Gaussian measure, called the free field measure. Let  $G$  be the green kernel associated to the Laplace-Beltrami operator on  $\Sigma$ . Let  $S$  be the generic element of  $\Sigma$ . We have, in this case:

$$(2) \quad E[x(S)x(S')] = G(S, S')$$

and  $G(S, S) = \infty$  in two dimension. This explains that we consider only smeared fields, and that the Gaussian measure of the free field lives on distributions in fact ([46]).

There is another way to construct non-linear random fields, parametrized by surfaces, which are continuous. It is the purpose of infinite dimensional processes on infinite dimensional manifolds, which comes from the work of Kuo ([23]), Beloposkaya-Daletskii ([8]) and Daletskii ([13]).

The simplest geometry for the world sheet is the case where we consider the cylinder: this corresponds to diffusion processes on the loop space of the manifold.

This can be seen in two ways:

-)Either we consider Ornstein-Uhlenbeck process on the loop space, by using Dirichlet forms theory ([6], [15], [27], [41]).

-)Or we consider Brownian motion on the loop space, in the sense of Airault-Malliavin ([3],[10], [11], [36]).

In these two cases, we say that we are in presence of a 1 + 1 dimensional field theory, 1 for the dynamical time and 1 for the dimension of the internal time of the theory.

Some refinement of 1 + 1 dimensional field theories were done in [12], [37] and [38], where we consider pants. It is a finite dimensional constrain, and traditional tools of stochastic analysis as quasi-sure analysis ([2], [48]) can be applied.

Until now, in order to understand a more complicated world-sheet, we were obliged to consider a 1 + 2 dimensional field theory, that is processes on the set of maps from  $\Sigma$  into  $M$  ([28], [30], [31], [32], [33], [34], [39]) and the associated heat-kernel measure.

Our goal is to understand a more complicated world-sheet, that is the sphere (We start from the constant loop and we arrive at the constant loop) as a 1 + 1 dimensional stochastic field theory. So we construct a process on the loop space of  $R^d$  (endowed with a non-trivial Riemannian structure, but **without cut-locus!**) starting from the constant loop and submitted to coming-up at the constant loop at time 1. It is an infinite dimensional constrain, and we cannot apply tools of quasi-sure analysis in order to do this work ([2], [48]).

We refer to the survey of Albeverio ([4]) for the relation between stochastic analysis and mathematical physics as well as the 3 surveys of Léandre about that ([27], [29], [40]).

## 2 The model

Let  $X_i$  be  $m + 1$  vector fields on  $R^d$  bounded with bounded derivatives of all orders. We suppose that the quadratic form  $g(x)^{-1}$

$$(3) \quad \xi \rightarrow \sum_{i=1}^m \langle X_i(x), \xi \rangle^2$$

is uniformly invertible on  $R^d$ . Let  $w^i, i = 1, \dots, m$  a  $R^m$ -valued Brownian motion. We can consider the Stratonovich differential equation issued from  $x$ :

$$(4) \quad dx_t(x) = \sum_{i=1}^m X_i(x_t(x))dw_t^i + X_0(x_t(x))dt$$

The law of  $x_t(x)$  has a smooth density  $q_t(x, y)$  with respect to the Lebesgue measure on  $R^d$ , strictly positive, and which is smooth in  $t > 0, x, y$ . If we constrain the diffusion  $x_t(x)$  to be  $x$  at time 1, it is the solution of the Stratonovich differential equation ([9], [45]):

$$(5) \quad dx_t(x, x) = \sum_{i=1}^m X_i(x_t(x, x))d\tilde{w}_t^i + X_0(x_t(x, x))dt$$

where

$$(6) \quad d\tilde{w}_t^i = dw_t^i + \langle X_i(x_t(x, x)), \text{grad} \log q_{1-t}(x_t(x, x), x) \rangle dt$$

We consider the Hilbert space  $H$  of free loops  $h$  in  $R^m$  endowed with the Hilbert structure:

$$(7) \quad \int_{S^1} |h(s)|^2 ds + \int_{S^1} |d/dsh(s)|^2 ds = \|h\|^2$$

Let  $\psi_s^i$  be the map  $h \rightarrow h^i(s)$  where  $h^i$  is the coordinate of order  $i$  of  $h$ . It is a continuous linear form. Therefore there exists an element  $e_s^i$  of  $H$  such that

$$(8) \quad h(s) = \langle h, e_s^i \rangle$$

Let us consider the Brownian motion  $B_t(\cdot)$  with values in  $H$ .  $t \rightarrow B_t(s)$  are a family of  $R^m$ -valued Brownian motion indexed by  $s \in S^1$ . Moreover,

$$(9) \quad E[B_t^i(s)B_t^j(s')] = t \langle e_s^i, e_{s'}^j \rangle$$

We deduce by Kolmogorov lemma ([44]), that the random field  $(t, s) \rightarrow B_t(s)$  is almost surely Hoelder. We consider the solution  $x_t(s, x)$  of (5) where we replace  $dw_t$  by  $dB_t(s)$  and the associated bridge  $x_t(s, x, x)$ .

We suppose in the sequel that the following hypothesis is satisfied:

**Hypothesis H:** For all multi-index  $(\alpha)$  of length equal to 2, we have for  $t \leq 1$ :

$$(10) \quad t(|\frac{\partial^{(\alpha)}}{\partial x^{(\alpha)}} \log q_t(x, y)| + |\frac{\partial^{(\alpha)}}{\partial y^{(\alpha)}} \log q_t(x, y)|) < C < \infty$$

We refer to [22] and to [42] about this hypothesis, which says more and less that the Riemannian manifold  $R^d$  endowed with the Riemannian metric  $g$  has no cut-locus.

Moreover, we can suppose that the diffusion  $x_t(s, x)$  is the Brownian motion on  $R^d$  endowed with the Riemannian structure  $g(x)$ , by choosing a suitable drift  $X_0$ .

### 3 The main theorem

**Theorem III.1:** the multiparameter process  $(t, s) \rightarrow x_t(s, x, x)$  has a version which is Hoelder in  $s$  and  $t$  under **Hypothesis H**.

**Lemma III.2:** When  $t \rightarrow 0$

$$(11) \quad P\{\sup_s |x_t(s, x, x) - x| > C\} \leq K' \exp[-K/t]$$

for  $K > 0$  depending only from  $C > 0$

**Proof:** We see that if we put

$$(12) \quad U_t(s, s') = x_t(s, x, x) - x_t(s', x, x)$$

that

$$(13) \quad U_t(s, s') = \int_0^t O_u^1 U_u(s, s') du + \int_0^t O_u^2 U_u(s, s') \delta B_u(s) + \int_0^t O_u^3 \delta(B_u(s) - B_u(s'))$$

where  $O_u^2$  and  $O_u^3$  are bounded continuous adapted processes and  $|O_u^1| \leq \frac{C}{1-u}$  by **Hypothesis H**.  $\delta$  denotes the Itô differential. We deduce by Hoelder inequality and Burkholder-Davies-Gundy inequality that for some  $k$  depending of  $p$  only:

$$(14) \quad E|U_t(s, s')|^p \leq C |\log(1-t)|^k \int_0^t \frac{du}{1-u} |U_u(s, s')|^p + C |s - s'|^{p/2}$$

because

$$(15) \quad E[|B_t(s) - B_t(s')|^2] = tO(1)|s - s'|$$

when  $s \rightarrow s'$

Let us recall that for arbitrary small  $\epsilon$

$$(16) \quad \int_0^t \frac{du}{(1-u)|\log(1-u)|^{k'}} \leq C(1-u)^{-\epsilon}$$

We deduce by Gronwall lemma that:

$$(17) \quad E[|U_t(s, s')|^p] \leq C |s - s'|^{p/2} \exp[K(1-t)^{-\epsilon}]$$

We deduce by Kolmogorov lemma ([44]) that  $s \rightarrow x_t(s, x, x)$  is  $1/2 - \epsilon'$  Hoelder with Hoelder norm in  $L^p$  bounded by  $\exp[K(1-t)^{-\epsilon}]$ .

We consider  $\exp[C_1(1-t)^{-1}]$  sites  $s_i$  on the circle. It is classical that  $P\{|x_t(s_i, x, x) - x| > C\}$  is bounded by  $\exp[-D(1-t)^{-1}]$  when  $t \rightarrow 0$ . We choose  $C_1$  small with respect of  $D$ . We get:

$$(18) \quad P\{\sup_s |x_t(s, x, x) - x| > C\} \leq \sum_i P\{|x_t(s_i, x, x) - x| > C\} + P\{\|x_t(\cdot, x, x)\|_{1/2-\epsilon'} > \exp[E(1-t)^{-1}]\}$$

for some convenient  $E > 0$  and where  $\|\cdot\|_{1/2-\epsilon'}$  denotes the Hoelder norm on the circle with exponent  $1/2 - \epsilon'$ .

We deduce the estimate

$$(19) \quad P\{\sup_s |x_t(s, x, x) - x| > C\} \leq \exp[C_1(1-t)^{-1}] \exp[-D(1-t)^{-1}] + \exp[-E(1-t)^{-1}] \exp[K(1-t)^{-\epsilon}]$$

The result because  $\epsilon$  is small and because  $C_1$  is strictly smaller than  $D$ .

◇

We consider a smooth function  $h$  from  $R$  into  $R^+$  equal to 0 in a neighborhood of 0 and equal to 1 in a neighborhood of  $\infty$ . We denote

$$(20) \quad z_t = h(\sup_s |x_t(s, x, x) - x|)$$

From Lemma III.2, we deduce:

**Lemma III.3:**  $Z = \int_0^1 \frac{z_t}{(1-t)} dt$  is almost surely finite.

Let us denote:

$$(21) \quad \sup_s |x_t(s, x, x) - x| \wedge C = z'_t$$

By a small modification of Lemma III. 2, we have if  $\alpha < 1/2$

$$(22) \quad P\{z'_t > (1-t)^\alpha\} \leq \exp[-C(1-t)^{-\epsilon}]$$

such that  $Z' = \int_0^1 \frac{z'_t}{(1-t)} dt$  is almost surely finite.

**Lemma III.4:** If  $Z + Z' < n$ , there exist a constant  $C(n, p)$  still denoted by  $C$  such that:

$$(23) \quad E[|x_t(s, x, x) - x_{t'}(s', x, x)|^p]^{1/p} \leq C(|s - s'|^{1/2} + |t - t'|^{1/2})$$

**Proof:** We look at (13). We write

$$(24) \quad O_u^1 = O_u^1 z_u + O_u^1(1 - z_u) = V_u^1 + V_u^2$$

We have  $|V_u^1| \leq C(1-u)^{-1} z_u$ . We work in in normal system of coordinates of  $x$  for the Riemannian structure such that when  $t \rightarrow 0$  we have the asymptotic expansion ([21], [26], [25], [50])

$$(25) \quad q_t(y, x) = t^{-d/2} \exp[-d^2(x, y)/2t] \sum c_i(y, x) t^i$$

if  $|x - y| < C$ . This asymptotic expansion works for the derivatives of  $q_t(x, y)$  too, by taking derivatives in  $x$  and  $y$  of the asymptotics expansion (This implies that  $c_i(x, y)$  is smooth in  $x$  and  $y$ ). We write in normal coordinates  $d^2(x, y) \sim |x - y|^2$  where  $d^2(x, y)$  is the Riemannian metric on  $R^d$  for the Riemannian structure  $g$ . We deduce that

$$(26) \quad V_u^2(U_t(s, s')) = -\frac{1}{1-t} U_t(s, s') + \frac{A_t}{1-t} U_t(s, s')$$

where

$$(27) \quad |A_t| \leq C(\sup_s |x_t(s, x, x)_x| \wedge C) = C z'_t$$

We deduce that  $U_t(s, s')$  satisfies

$$(28) \quad \delta U_t(s, s') = -\frac{1}{1-t}U_t(s, s')dt + \frac{A_t}{1-t}U_t(s, s')dt + V_t^1 U_t(s, s')dt + O_t^2 U_t(s, s')\delta B_t(s) + O_t^3 \delta(B_t(s) - B_t(s'))$$

We put  $U_t(s, s') = (1-t)C_t(s, s')$ .  $C_0(s, s') = 0$  such that  $C_t(s, s')$  satisfies the differential equation:

$$(29) \quad \delta C_t(s, s') = \frac{A_t}{(1-t)}C_t(s, s')dt + V_t^1 C_t(s, s')dt + O_t^2 C_t(s, s')\delta B_t(s) + \frac{O_t^3}{(1-t)}\delta(B_t(s) - B_t(s'))$$

We deduce by Gronwall inequality since  $Z + Z' \leq n$  that

$$(30) \quad |C_t(s, s')| \leq C'_n \sup_{v \leq t} (|\int_0^v O_u^2 C_u(s, s')\delta B_t(s) + \frac{O_u^3}{1-u}\delta(B_u(s) - B_u(s'))|)$$

By Gronwall lemma and Burkholder -Davis-Gundy inequality, we deduce that

$$(31) \quad \|C_t(s, s')\|_{L^p} \leq C''_n |s - s'|^{1/2} (1-t)^{-1}$$

Therefore

$$(32) \quad \|U_t(s, s')\|_{L^p} \leq C'''_n |s - s'|^{1/2}$$

Moreover, we have clearly the inequality:

$$(33) \quad \|x_t(s, x, x) - x_{t'}(s, x, x)\|_{L^p} \leq C|t - t'|^{1/2}$$

◇

**Proof of Theorem III.1:** If  $Z + Z' < n$ , we have an Hoelder version of the random field  $(t, s) \rightarrow x_t(s, x, x)$ . Since  $Z + Z'$  is almost surely finite the result arises.

◇

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